Mu2e: coherent $\mu \rightarrow e$ conversion experiment at Fermilab

MU2e

Gianantonio Pezzullo INFN of Pisa

on behalf of the Mu2e collaboration

MUSE The Mu2e collaboration





Argonne National Laboratory, Boston University, Brookhaven National Laboratory University of California, Berkeley, University of California, Irvine, California Institute of Technology, City University of New York, Joint Institute for Nuclear Research, Dubna, Duke University, Fermi National Accelerator Laboratory, Laboratori Nazionali di Frascati, Helmholtz-Zentrum Dresden-Rossendorf, University of Houston, University of Illinois, INFN Genova, Kansas State University, Lawrence Berkeley National Laboratory, INFN Lecce and Università del Salento, Lewis University, University of Louisville, Laboratori Nazionali di Frascati and Università Marconi Roma, University of Minnesota, Muons Inc., Northern Illinois University, Northwestern University, Novosibirsk State University/Budker Institute of Nuclear Physics, Institute for Nuclear Research, Moscow, INFN Pisa, Purdue University, Rice University, University of South Alabama, Sun Yat Sen University, University of Virginia, University of Washington, Yale University

MUSE What is $\mu \rightarrow e$ conversion

- μ converts to an electron in the presence of a nucleus $\mu^- N
 ightarrow e^- N$

$$E_e = m_{\mu} c^2 - B_{\mu}(Z) - C(A) = 104.973 MeV$$

- for Aluminum: $\begin{cases} B_{\mu}(Z) \text{ is the muon binding energy (0.48 MeV)} \\ C(A) \text{ is the nuclear recoil energy (0.21 MeV)} \end{cases}$
- µ conversion in the SM is introduced by the neutrino masses and mixing at a negligible level ~ 10⁻⁵²
- Many SM extensions enhance the rate through mixing in the high energy sector of the theory (other particles in the loop...)







• Mu2e will improve by a factor 10⁴ the present best limit!

Experimental setup



• Production Solenoid:

- Proton beam strikes target, producing mostly pions
- Graded magnetic field contains backwards pions/muons and reflects slow forward pions/muons

• Detector Solenoid:

- ➡ Capture muons on Al target
- Measure momentum in tracker and energy in calorimeter
- ➡ Graded field "reflects" downstream conversion electrons emitted upstream



- Transport Solenoid:
- Select low momentum, negative muons
- ➡ Antiproton absorber in the mid-section







• µ decay-in-orbit

- Cosmic-induced background
- Antiproton-induced background
- Radiative π capture

MUSE μ decay-in-orbit (DIO)



Czarnecki et al., arXiv:1106.4756v2 [hep-ph] Phys. Rev. D 84, 013006 (2011)

G. Pezzullo (INFN of Pisa)

NFN





• µ decay-in-orbit:

✓ low-mass tracker with high performance

- Cosmic-induced background
- Antiproton-induced background
- Radiative π capture





- µ decay-in-orbit:
 - \checkmark low-mass tracker with high performance
- Cosmic-induced background:

 \checkmark cosmic ray veto and PID

- Antiproton-induced background
- Radiative π capture



Cosmic Ray Veto



- Veto system covers entire DS and half TS
- 4 layers of scintillator
 - each bar is 5x2x~450 cm³
 - ✤ 2 WLS fibers/bar
 - read out at both ends with SiPM
- inefficiency < 10^{-4}







MUSE Physics background



- µ decay-in-orbit:
 - \checkmark low-mass tracker with high performance
- Cosmic-induced background:
 - \checkmark cosmic ray veto and PID
- Antiproton-induced background
 - ✓ beam line and PID
- Radiative π capture



Antiproton absorber



- p-bar reaching the Detector
 Solenoid can annihilate in the Al stopping target
 ~ 2GeV of shower created
- 2 Be absorbers in the Transport Solenoid (TS) are used to limit the p-bar flux:
 - ✤ at the entrance of the TS
 - in the middle of the TS

Transport solenoid - top view



MUSE Physics background



• µ decay-in-orbit:

 \checkmark low-mass tracker with high performance

• Cosmic-induced background:

 \checkmark cosmic ray veto and PID

Antiproton-induced background

 \checkmark beam line and PID

• Radiative π capture: $\pi^- + N \rightarrow \gamma + N'$

✓ pulsed beam and extinction of out-of-time protons



Pulsed beam



- Beam period : 1.7 μ s ~ 2 x τ_{μ}^{Al}
- Beam intensity: 3.15×10^7 p/bunch
- duty cycle :~ 30%
- out-of-time protons / in-time protons < 10^{-10}



MUSE-





- 18 stations with straws transverse to the beam
- Straw technology employed:
 ✓ 5 mm diameter, 15 µm Mylar walls
 ✓ 25 µm Au-plated W sense wire
 ✓ 80/20 Ar/CO₂ with HV ~ 1500 V
- Inner 38 cm un-instrumented:
 - \checkmark blind to beam flash
 - ✓ blind to low pT charged particles coming from the Al target













Calorimeter



- 2 disks; each disk contains 860 undoped CsI crystals $20 \times 3.4 \times 3.4$ cm³
- Readout by large area MPPC + waveform digitizer boards @ 200 MHz
- Allows to measure: E/p and TOF to provide Particle identification
- Improve track search via a calorimeter-seed pattern recognition
- Time resolution σ_t < 200 ps @ 100 MeV measured @ BTF in Frascati



Calorimeter prototype







Beam test @ BTF in Frascati

28th Rencontres de Blois - June I 2016

G. Pezzullo (INFN of Pisa)

MUSE Mu2e signal sensitivity

Reconstructed e Momentum



• Single-event-sensitivity = 2.9×10^{-17}

• Total background < 0.5 events



Conclusions



- \bullet Mu2e is an experiment to search for CLFV in μ coherent conversion
 - \checkmark aims 4 orders of magnitude improvement

 \checkmark expected SES = 2.9 x 10⁻¹⁷

- \checkmark any signal would be an un-ambiguous proof of physics BSM
- R&D phase mostly completed
- Civil construction and magnets procurement already started
- Data taking starts on 2021





backup slides



MUSE- Muonic atom life times



G. Pezzullo (INFN of Pisa)

NFN



 $R_{\mu e}$ rate vs Z







Mu2e signal?





- A next-generation Mu2e experiment makes sense in all scenarios:
 - \checkmark Push sensitivity or
 - \checkmark Study underlying new physics
 - ✓ Will need more protons upgrade accelerator
 - ✓ Snowmass white paper, arXiv:1307.116

MUSE- Model independent Lagrangian



G. Pezzullo (INFN of Pisa)

28th Rencontres de Blois - June I 2016



CLFV limits I



Process	Upper limit		
$\mu^+ \to e^+ \gamma$	$< 5.7 \times 10^{-13}$		
$\mu^+ \to e^+ e^- e^+$	$< 1.0 \times 10^{-12}$		
$\mu^{-}\mathrm{Ti} \rightarrow e^{-}\mathrm{Ti}$	$< 1.7 \times 10^{-12}$		
$\mu^{-}\mathrm{Au} \to e^{-}\mathrm{Au}$	$< 7 \times 10^{-13}$		
$\mu^+ e^- \to \mu^- e^+$	$< 3.0 \times 10^{-13}$		
$\tau \to e\gamma$	$< 3.3 \times 10^{-8}$		
$\tau^- \to \mu \gamma$	$< 4.4 \times 10^{-8}$		
$\tau^- \to e^- e^+ e^-$	$< 2.7 \times 10^{-8}$		
$\tau^- \to \mu^- \mu^+ \mu^-$	$< 2.1 \times 10^{-8}$		
$\tau^- \to e^- \mu^+ \mu^-$	$< 2.7 \times 10^{-8}$		
$\tau^- \to \mu^- e^+ e^-$	$< 1.8 \times 10^{-8}$		
$\tau^- \to e^+ \mu^- \mu^-$	$< 1.7 \times 10^{-8}$		
$\tau^- \to \mu^+ e^- e^-$	$< 1.5 \times 10^{-8}$		

G. Pezzullo (INFN of Pisa)



CLFV limits 2



Process	Upper limit
$\pi^0 \to \mu e$	$< 8.6 \times 10^{-9}$
$\mathrm{K}^{0}_{\mathrm{L}} \to \mu e$	$< 4.7 \times 10^{-12}$
$K^+ \to \pi^+ \mu^+ e^-$	$< 2.1 \times 10^{-10}$
$\mathrm{K}^{0}_{\mathrm{L}} \to \pi^{0} \mu^{+} e^{-}$	$< 4.4 \times 10^{-10}$
$Z^0 \to \mu e$	$< 1.7 \times 10^{-6}$
$Z^0 \to \tau e$	$< 9.8 \times 10^{-6}$
$Z^0 \to \tau \mu$	$< 1.2 \times 10^{-6}$



Out-of-time protons



- The RF structure of the Recycler provides some "intrinsic" extinction:
 ✓ Intrinsic extinction ~10⁻⁵
- A custom-made AC dipole placed just upstream of the production solenoid provides additional extinction:
 ✓ AC dipole extinction ~ 10⁻⁶ - 10⁻⁷
- Together they provide a total extinction:
 ✓ Total extinction ~ 10⁻¹¹ 10⁻¹²
- Extinction measured using a detector system: Si-pixel + sampling EMC







phase I

phase II





Why Particle Identification is needed



- **Cosmic ray** and **antiproton** induced background can be divided into 2 main categories:
 - I. e⁻ generated via interactions producing a track mimicking the CE
 - 2. non-electron particles (μ and π) that are reconstructed as an "electron-like" track mimicking the CE
- (1) represents the irreducible background, while (2) can be suppressed using a PID method

Mu2e PID method:

- Information from reconstructed tracks and calorimeter clusters are combined for identifying group (2)
- Stringent requirement from Cosmic: μ -rejection factor ≥ 200

ISE Cosmic µ rejection

- 105 MeV/c e⁻ are ultra-relativistic, while 105 MeV/c μ have β ~ 0.7 and a kinetic energy of ~ 40 MeV
- Likelihood rejection combines $\Delta t = t_{track} t_{cluster}$ and E/p:

$$\ln L_{e,\mu} = \ln P_{e,\mu}(\Delta t) + \ln P_{e,\mu}(E/p)$$



PID performance



• A muon-rejection of 200 corresponds to a cut at $\ln L_{e/\mu} > 1.5$ and an e⁻ efficiency of ~ 96%

G. Pezzullo (INFN of Pisa)



PID yield vs performance





• In the range $\sigma_E/E<0.1$ and $\sigma_t<0.5$ ns the e⁻ efficiency is within 2%





Crystals properties



Crystal	\mathbf{BaF}_2	LYSO	CsI
Density $[g/cm^3]$	4.89	7.28	4.51
Radiation length [cm] X_0	2.03	1.14	1.86
Molière radius [cm] R _m	3.10	2.07	3.57
Interaction length [cm]	30.7	20.9	39.3
dE/dx [MeV/cm]	6.5	10.0	5.56
Refractive Index at λ_{\max}	1.50	1.82	1.95
Peak luminescence [nm]	220,300	402	310
Decay time τ [ns]	0.9,650	40	16
Light yield (compared to $NaI(TI)$) [%]	4.1, 3.6	85	3.6
Light yield variation with	0.1, -1.9	-0.2	-1.4
temperature $[\%/^{\circ}C]$			
Hygroscopicity	Slight	None	Slight



Csl properties





$$\mathrm{EWLT} = \frac{\int \mathrm{LT}(\lambda) \mathrm{Em}(\lambda) d\lambda}{\int \mathrm{Em}(\lambda) d\lambda}$$

• where $LT(\lambda)$ is the light transmittance and $Em(\lambda)$ is the emission spectrum





Csl Light Output





G. Pezzullo (INFN of Pisa)



CsI LRU





Csl rad damage













Energy resolution



- Prototype dimensions: I.3 R_{Moliere}² x I0 X₀
- Still comparison between data and Monte Carlo useful







- Average tof from middle of the tracker to the calorimeter ~ 8 ns
- Mean drift time ~ 20 ns
- Difference of these two numbers is consistent with the peak position

G. Pezzullo (INFN of Pisa)

MUSE Cluster positon selection





- Graded magnetic field between the stopping target and the tracker limits the CE pT
- Cluster position identifies the semi-plane where the CE track relies

G. Pezzullo (INFN of Pisa)

MUSEMu2e track reconstruction



The Mu2e track reconstruction has several specific features:

- a CE makes 2-3 full turns in the tracker
- time dependence of the track-hit position:

 $r_{\text{drift}} = v_{\text{drift}} \cdot (t_{\text{measured}} - T_0 - t_{\text{flight}})$

The track reconstruction is factorized into 2 main steps:

- I. Track finding: provides a set of straw hits consistent with a track candidate
- 2. Kalman based fitter: performs the final reconstruction

The **track finding** uses two algorithms:

- A. Standalone: relies only on the tracker information
- B. Calorimeter-seeded: seeds the track search using the reco cluster