

The Mu2e collaboration





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University of California, Berkeley, University of California, Irvine, California Institute of Technology, City University of New York,
Joint Institute for Nuclear Research, Dubna, Duke University, Fermi National Accelerator Laboratory, Laboratori Nazionali di
Frascati, Helmholtz-Zentrum Dresden-Rossendorf, University of Houston, University of Illinois, INFN Genova, Kansas State
University, Lawrence Berkeley National Laboratory, INFN Lecce and Università del Salento, Lewis University, University of Louisville,
Laboratori Nazionali di Frascati and Università Marconi Roma, University of Minnesota, Muons Inc., Northern Illinois University,
Northwestern University, Novosibirsk State University/Budker Institute of Nuclear Physics, Institute for Nuclear Research, Moscow,
INFN Pisa, Purdue University, Rice University, University of South Alabama, Sun Yat Sen University, University of Virginia, University
of Washington, Yale University

What is µ→e conversion



ullet μ converts to an electron in the presence of a nucleus $\,\mu^- N o e^- N$

$$E_e = m_{\mu} c^2 - B_{\mu}(Z) - C(A) = 104.973 \text{ MeV}$$

- for Aluminum: $\begin{cases} B_{\mu}(Z) \text{ is the muon binding energy } (0.48 \text{ MeV}) \\ C(A) \text{ is the nuclear recoil energy } (0.21 \text{ MeV}) \end{cases}$
- Signal normalization:

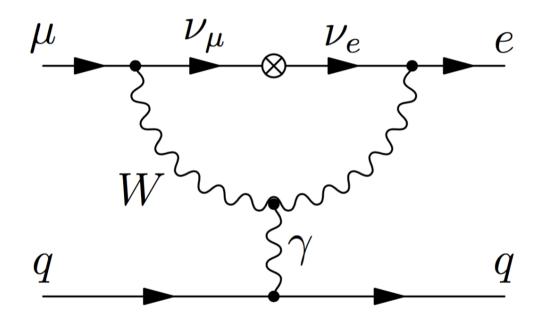
$$R_{\mu e} = \frac{\Gamma(\mu^{-} + N \to e^{-} + N)}{\Gamma(\mu^{-} + N \to all \text{ captures})}$$



CLFV in the SM



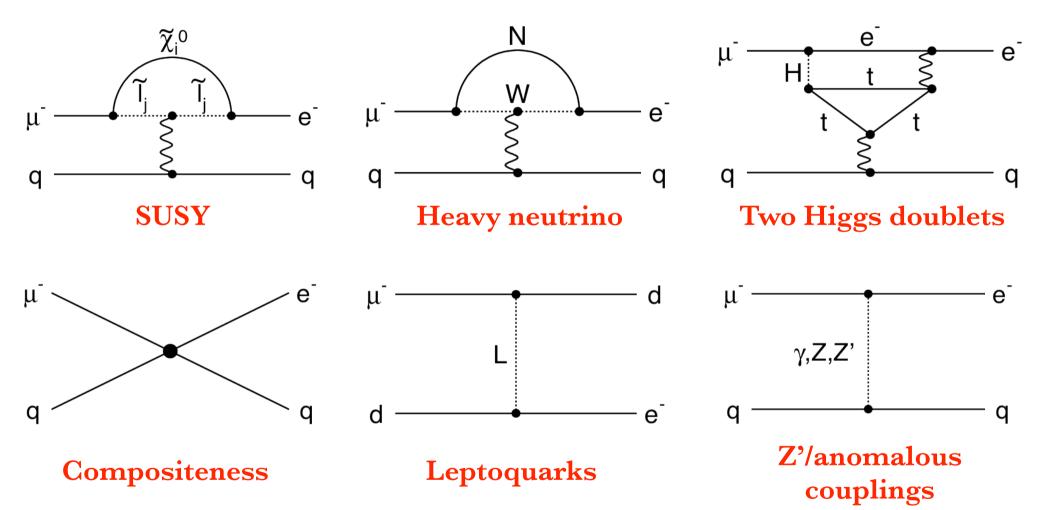
- CLFV process forbidden in the SM
- µ conversion in the SM is introduced by the neutrino masses and mixing at a negligible level ~ 10⁻⁵²



 Many SM extensions enhance the rate through mixing in the high energy sector of the theory (other particles in the loop...)

MUSE NP contributions to µ→e



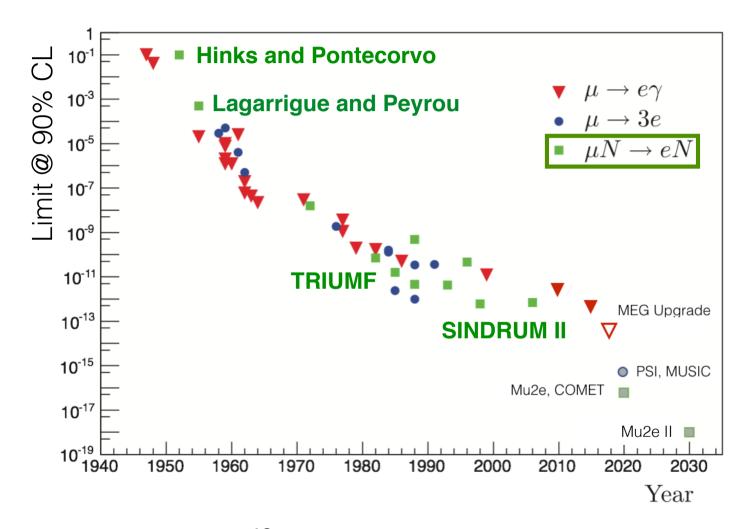


Any signal observation would be an unambiguous sign of NP



Historical perspective





- Best limit: $R_{\mu e} < 7 \times 10^{-13}$ by SINDRUM II at PSI [Eur.Phys.J C47(2006)]
- Mu2e will improve by a factor 10⁴

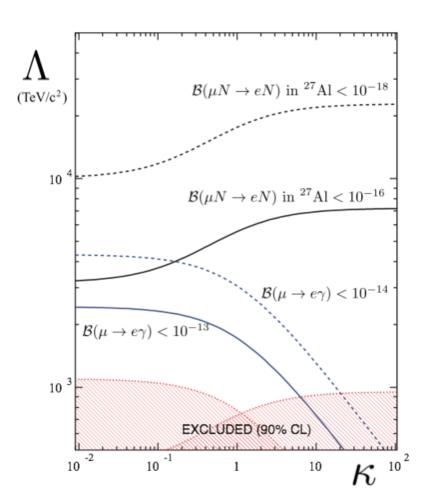
MUSE Model independent Lagrangian

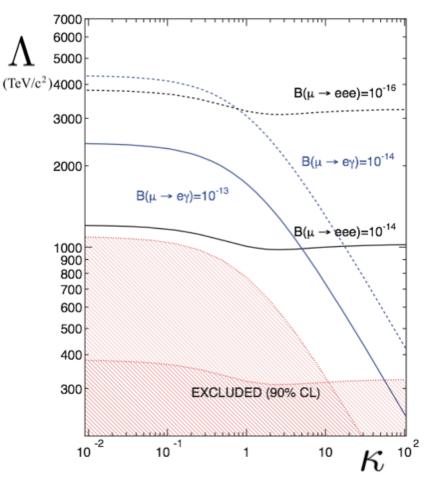


$$L_{\rm CLFV} = \frac{m_{\mu}}{(\kappa + 1)\Lambda^2} \bar{\mu}_R \sigma_{\mu\nu} e_L F^{\mu\nu} + \frac{\kappa}{(\kappa + 1)\Lambda^2} \bar{\mu}_L \gamma_{\mu} e_L \left(\bar{e}\gamma^{\mu}e\right)$$

"dipole term"

"contact term"





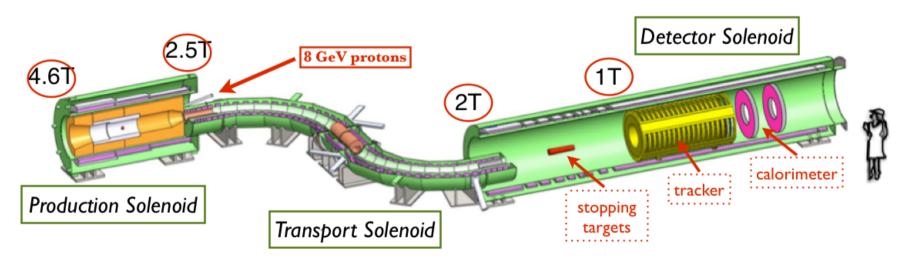


Experimental setup



- Production Solenoid:
- → Proton beam strikes target, producing mostly pions
- → Graded magnetic field contains backwards pions/muons and reflects slow forward pions/muons

- Detector Solenoid:
- → Capture muons on Al target
- → Measure momentum in tracker and energy in calorimeter
- → Graded field "reflects" downstream conversion electrons emitted upstream



- Transport Solenoid:
- → Select low momentum, negative muons
- → Antiproton absorber in the mid-section



Mu2e Detector

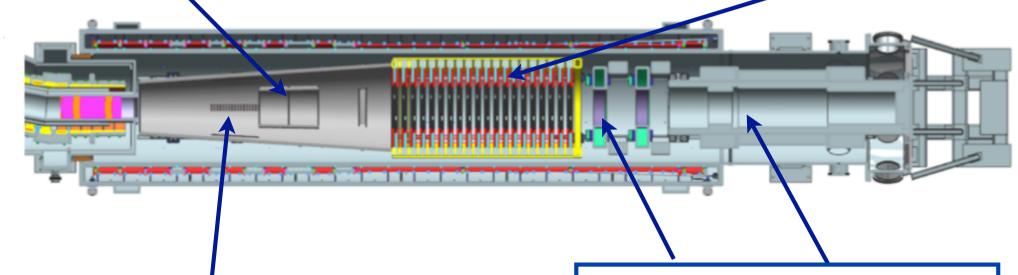


Proton absorber:

- ❖ made of high-density polyethylene
- ❖ designed in order to reduce proton flux on the tracker and minimize energy loss

• Tracker:

- ❖ ~20k straw tubes arranged in planes on stations, the tracker has 18 stations
- ❖ Expected momentum resolution < 200 keV/c



• Targets:

❖ 34 Al foils; Aluminum was selected mainly for the muon lifetime in capture events (**864 ns**) that matches nicely the need of prompt separation in the Mu2e beam structure.

• Calorimeter:

❖ 2 disks composed of undoped CsI crystals

Muon beam stop:

* made of several cylinders of different materials: stainless steel, lead and high density polyethylene



Muonic atoms



- Stopped μ^- is captured in atomic orbits
 - quickly (~fs) cascades into IS state
- Bohr radius ~20 fm (for Al)
 - \Rightarrow significant overlap between the μ^- and nucleus wave-functions
- For a μ^- in orbit three processes may happen:
 - ightharpoonup decay (39%): $\mu^- N o e^- ar{
 u}_e
 u_\mu N$, background
 - ightharpoonup capture (61%): $\mu^- + N o
 u_\mu + N'$, normalization
 - ightharpoonup conversion (<10⁻¹³): $\mu^- + N o e^- + N$, signal



Mu2e Detector

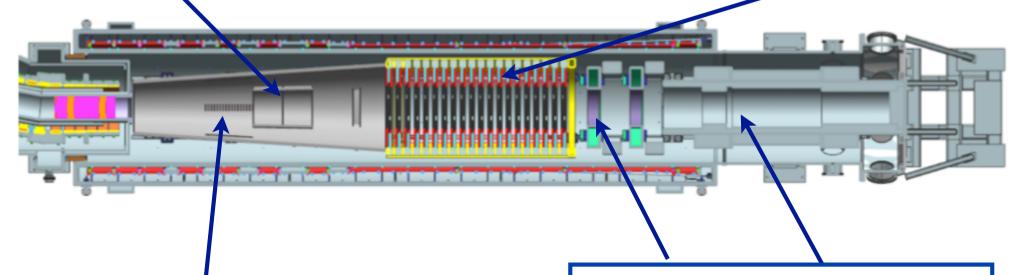


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Physics background

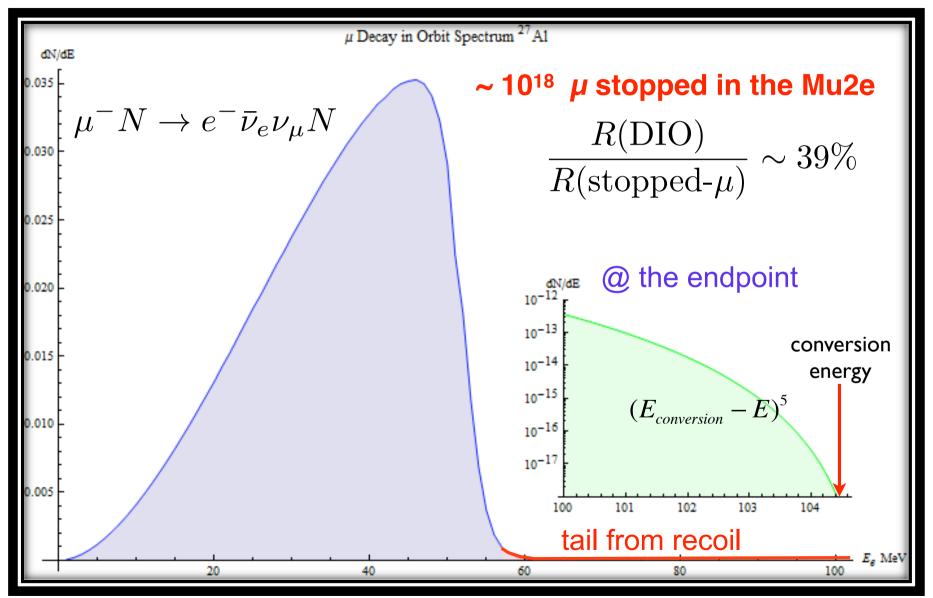


- μ decay-in-orbit
- Cosmic-induced background
- Antiproton-induced background
- Radiative π capture



μ decay-in-orbit (DIO)





Czarnecki et al., arXiv:1106.4756v2 [hep-ph] Phys. Rev. D 84, 013006 (2011)

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Physics background



- μ decay-in-orbit:
 - ✓ low-mass tracker with high performance
- Cosmic-induced background
- Antiproton-induced background
- Radiative π capture



Physics background



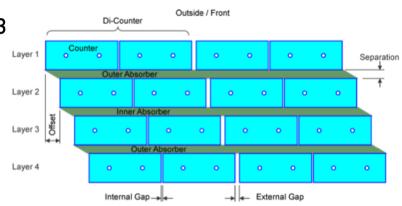
- μ decay-in-orbit:
 - √ low-mass tracker with high performance
- Cosmic-induced background:
 - √ cosmic ray veto and PID
- Antiproton-induced background
- Radiative π capture



Cosmic Ray Veto



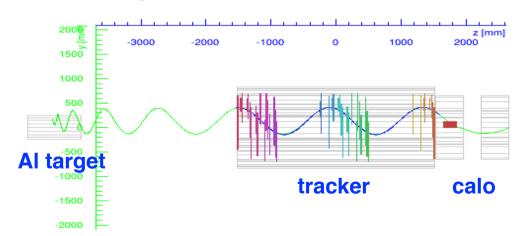
- Veto system covers entire DS and half TS
- 4 layers of scintillator
 - ♦ each bar is 5x2x~450 cm³
 - ❖ 2 WLS fibers/bar
 - read out with SiPM
- inefficiency < 10⁻⁴





TS-hole Cryo-hole

μ mimicking the signal





Physics background



- μ decay-in-orbit:
 - √ low-mass tracker with high performance
- Cosmic-induced background:
 - √ cosmic ray veto and PID
- Antiproton-induced background
 - ✓ beam line and PID
- Radiative π capture

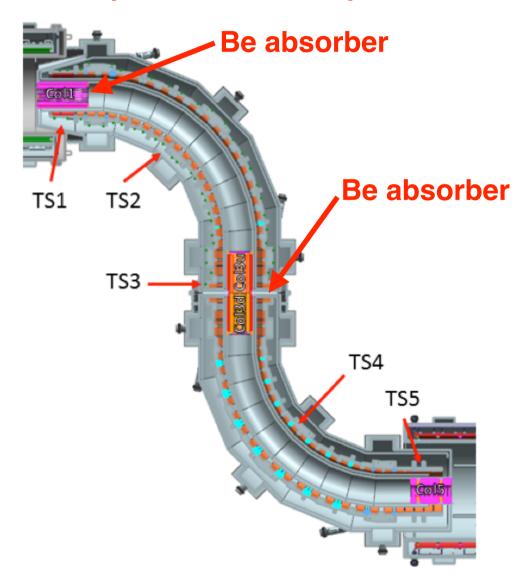


Antiproton absorber



- p-bar reaching the Detector
 Solenoid can annihilate in the Al stopping target
 - ❖ ~ 2GeV of shower created
- 2 Be absorbers in the Transport Solenoid (TS) are used to limit the p-bar flux:
 - * at the entrance of the TS
 - ❖ in the **middle** of the TS

Transport solenoid - top view





Physics background



- μ decay-in-orbit:
 - √ low-mass tracker with high performance
- Cosmic-induced background:
 - √ cosmic ray veto and PID
- Antiproton-induced background
 - √ beam line and PID
- Radiative π capture: $\pi^- + N \rightarrow \gamma + N'$
 - ✓ pulsed beam and extinction of out-of-time protons



Pulsed beam



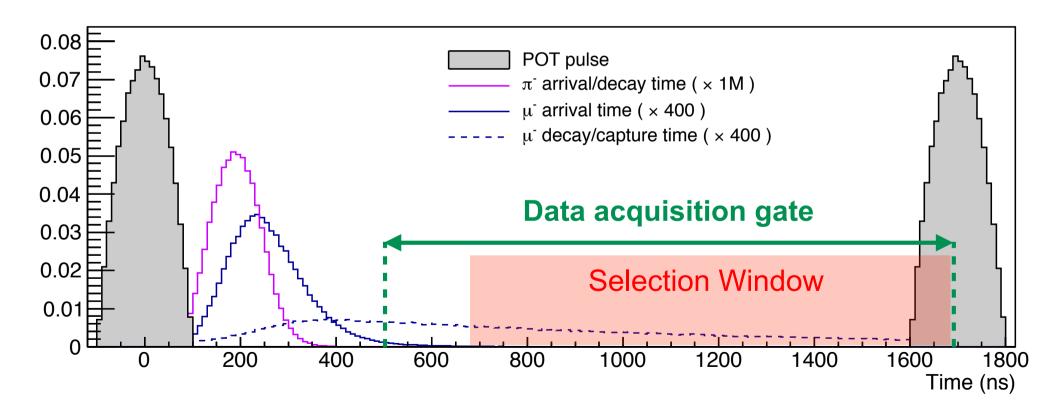
• Beam period : I.7 μ s ~ 2 x τ_{μ}^{Al}

• Beam intensity: 3.15×10^7 p/bunch

• duty cycle : ~ 30%

out-of-time protons / in-time protons < 10⁻¹⁰







Tracker

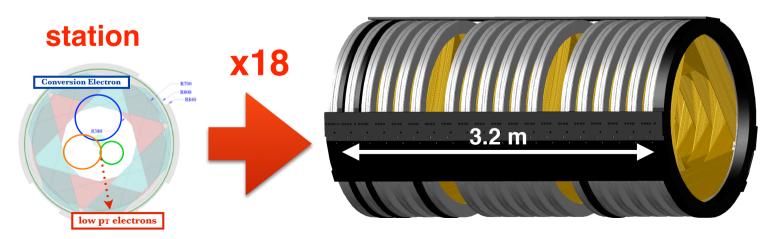


- 18 stations with straws transverse to the beam
- Straw technology employed:
 - √ 5 mm diameter, I5 µm Mylar walls
 - ✓ 25 µm Au-plated W sense wire
 - √ 80/20 Ar/CO₂ with HV ~ I500 V
- Inner 38 cm un-instrumented:
 - ✓ blind to beam flash
 - ✓ blind to low pT charged particles coming from the Al target
- Expected $\sigma_p < 200 \text{ keV/c}$











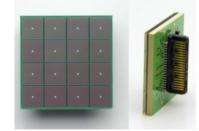


Calorimeter



- 2 disks; each disk contains 860 undoped CsI crystals 20 x 3.4 x 3.4 cm³
- Readout by large area MPPC + waveform digitizer boards @ 200 MHz
- Allows to measure: E/p and TOF to provide Particle identification
- Improve track search via a calorimeter-seed pattern recognition
- Time resolution σ_t < 200 ps @ 100 MeV measured @ BTF in Frascati

undoped Csl



MPPC

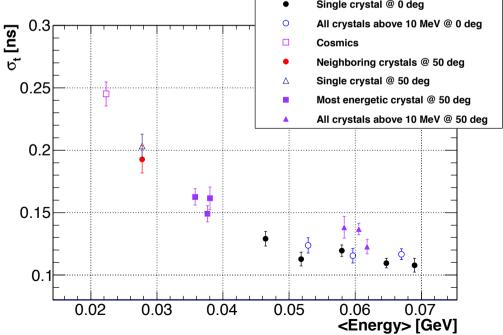










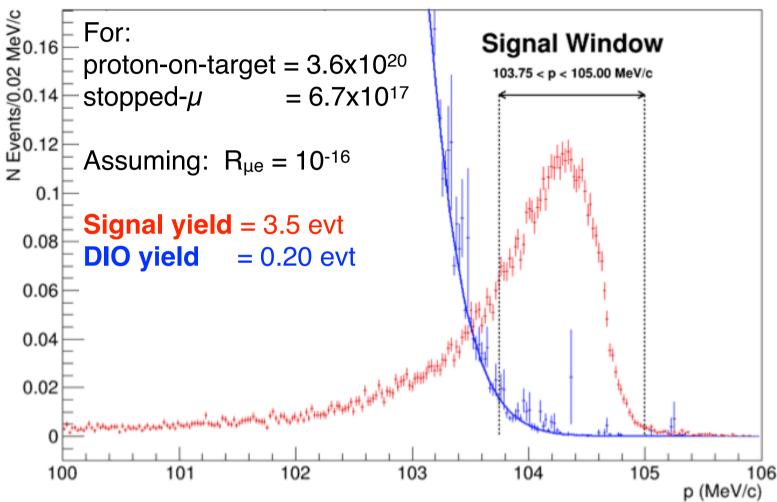




Mu2e signal sensitivity



Reconstructed e Momentum



- Single-event-sensitivity = 2.9×10^{-17}
- Total background < 0.5 events



Conclusions



- Mu2e is an experiment to search for CLFV in μ coherent conversion
 - ✓ aims 4 orders of magnitude improvement
 - ✓ unprecedented sensitivity to NP with mass scale up to ~ 10⁴ TeV
- Civil construction and magnets procurement already started
- R&D mature with data taking scheduled on 2021
- More info: http://mu2e.fnal.gov



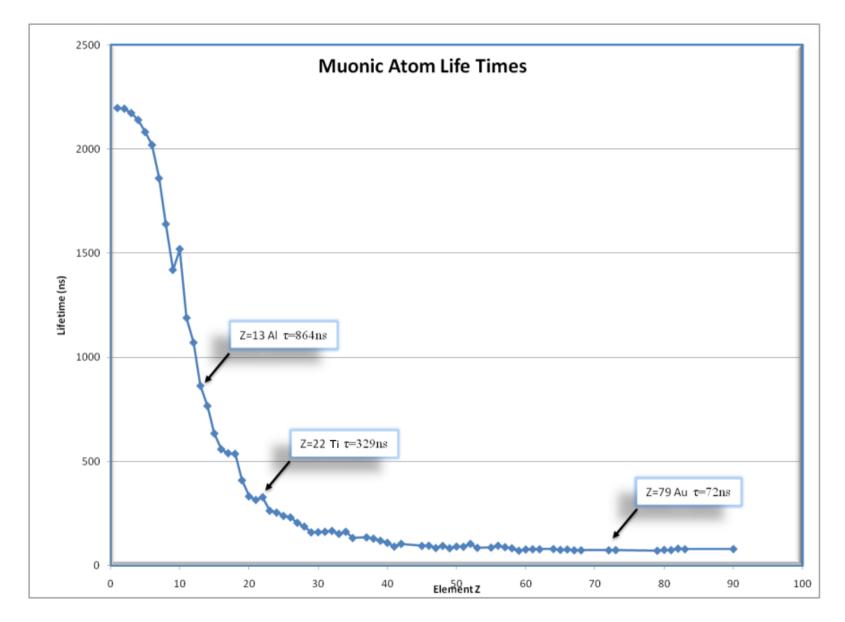






Muonic atom life times

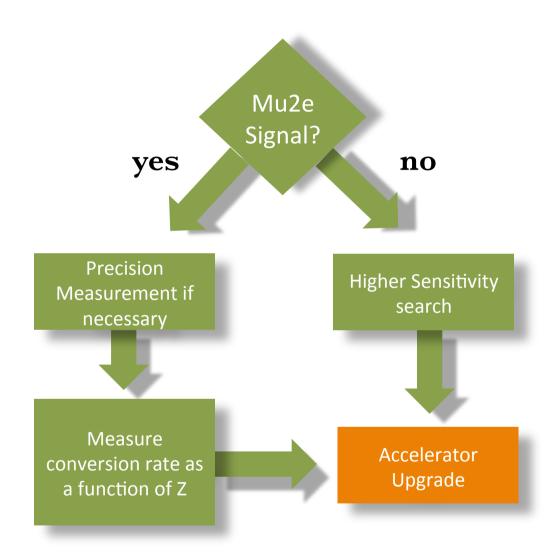






Mu2e signal?



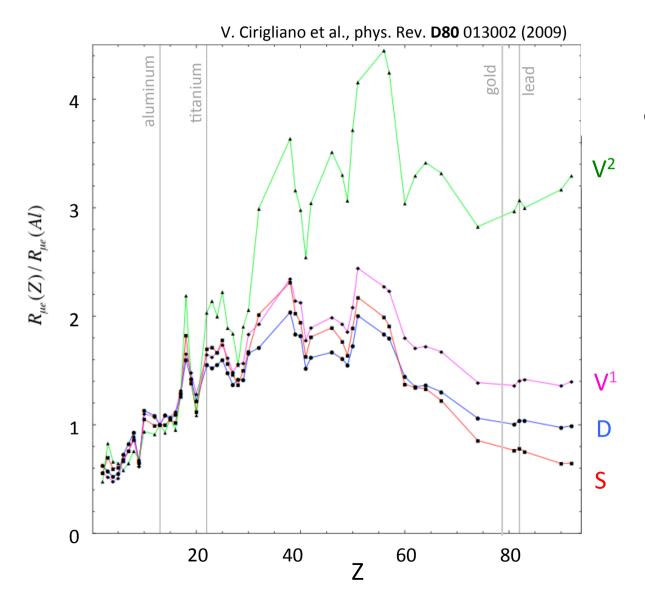


- A next-generation Mu2e experiment makes sense in all scenarios:
 - √ Push sensitivity or
 - √ Study underlying new physics
- Proton driver upgrade
 - https://pip2.fnal.gov
- Snowmass report:
 - √ arXiv:1311.5278



R_{µe} rate vs Z



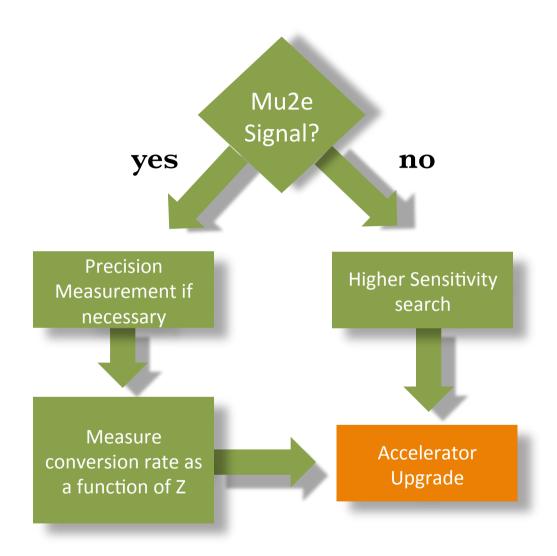


 Can use ratio of rates to determine dominant operator contribution



Mu2e signal?



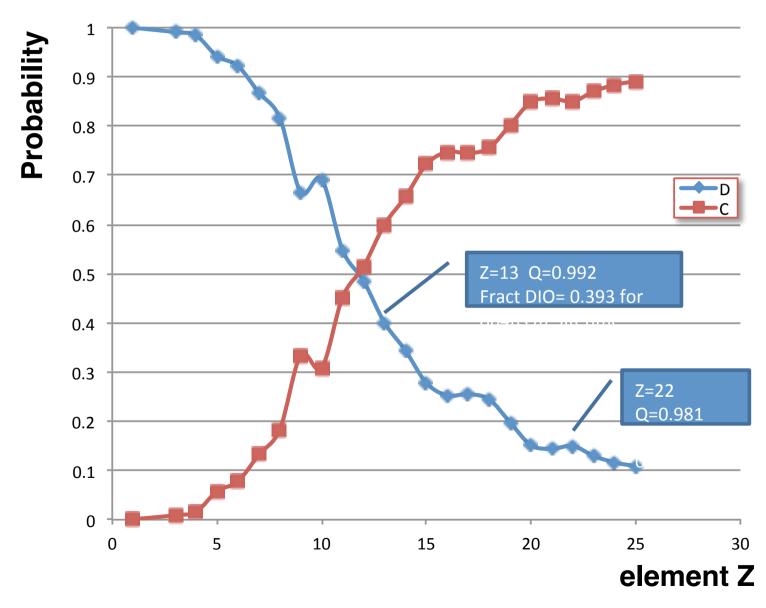


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MUSE Decay/Capture rate vs Z







CLFV limits I



Process	Upper limit
$\mu^+ \to e^+ \gamma$	$< 5.7 \times 10^{-13}$
$\mu^+ \to e^+ e^- e^+$	$< 1.0 \times 10^{-12}$
$\mu^- \mathrm{Ti} \to e^- \mathrm{Ti}$	$< 1.7 \times 10^{-12}$
$\mu^{-}\mathrm{Au} \to e^{-}\mathrm{Au}$	$< 7 \times 10^{-13}$
$\mu^+e^- \to \mu^-e^+$	$< 3.0 \times 10^{-13}$
$\tau \to e \gamma$	$< 3.3 \times 10^{-8}$
$\tau^- \to \mu \gamma$	$< 4.4 \times 10^{-8}$
$\tau^- \rightarrow e^- e^+ e^-$	$< 2.7 \times 10^{-8}$
$\tau^- \to \mu^- \mu^+ \mu^-$	$< 2.1 \times 10^{-8}$
$\tau^- \to e^- \mu^+ \mu^-$	$< 2.7 \times 10^{-8}$
$\tau^- \to \mu^- e^+ e^-$	$< 1.8 \times 10^{-8}$
$\tau^- \to e^+ \mu^- \mu^-$	$< 1.7 \times 10^{-8}$
$\tau^- \to \mu^+ e^- e^-$	$< 1.5 \times 10^{-8}$



CLFV limits 2

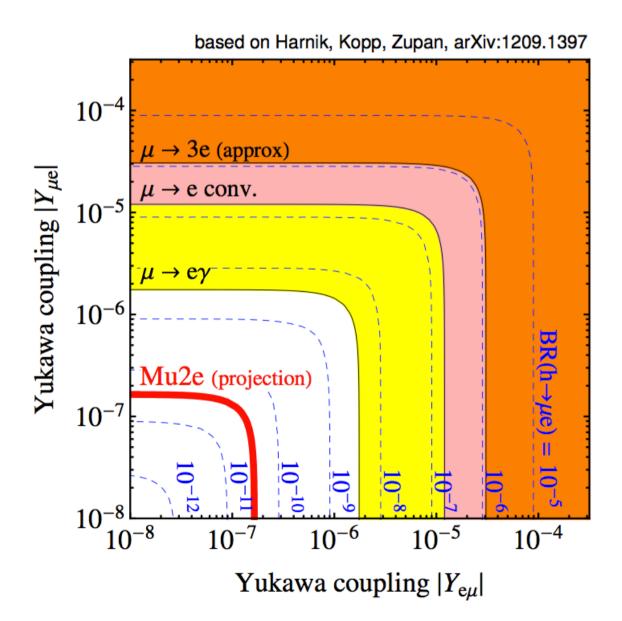


Process	Upper limit
$\pi^0 \to \mu e$	$< 8.6 \times 10^{-9}$
$\mathrm{K_L^0} \to \mu e$	$< 4.7 \times 10^{-12}$
$K^+ \to \pi^+ \mu^+ e^-$	$< 2.1 \times 10^{-10}$
$K_L^0 \to \pi^0 \mu^+ e^-$	$< 4.4 \times 10^{-10}$
$Z^0 \to \mu e$	$< 1.7 \times 10^{-6}$
$Z^0 \to \tau e$	$< 9.8 \times 10^{-6}$
$Z^0 \to \tau \mu$	$< 1.2 \times 10^{-6}$



Higgs couplings to µe





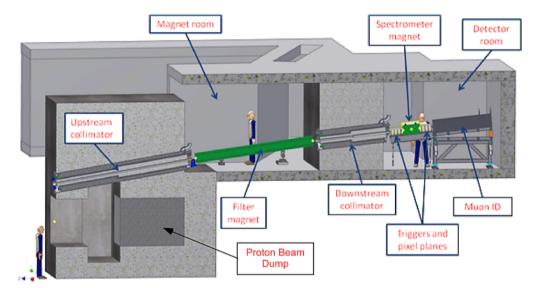
- Mu2e explores BR down to 10-11
- Outside the LHC reach



Out-of-time protons



- The RF structure of the Recycler provides some "intrinsic" extinction:
 - ✓ Intrinsic extinction ~10⁻⁵
- A custom-made AC dipole placed just upstream of the production solenoid provides additional extinction:
 - **✓** AC dipole extinction $\sim 10^{-6}$ 10^{-7}
- Together they provide a total extinction:
 - ✓ Total extinction $\sim 10^{-11}$ 10^{-12}
- Extinction measured using a detector system: Si-pixel + sampling EMC



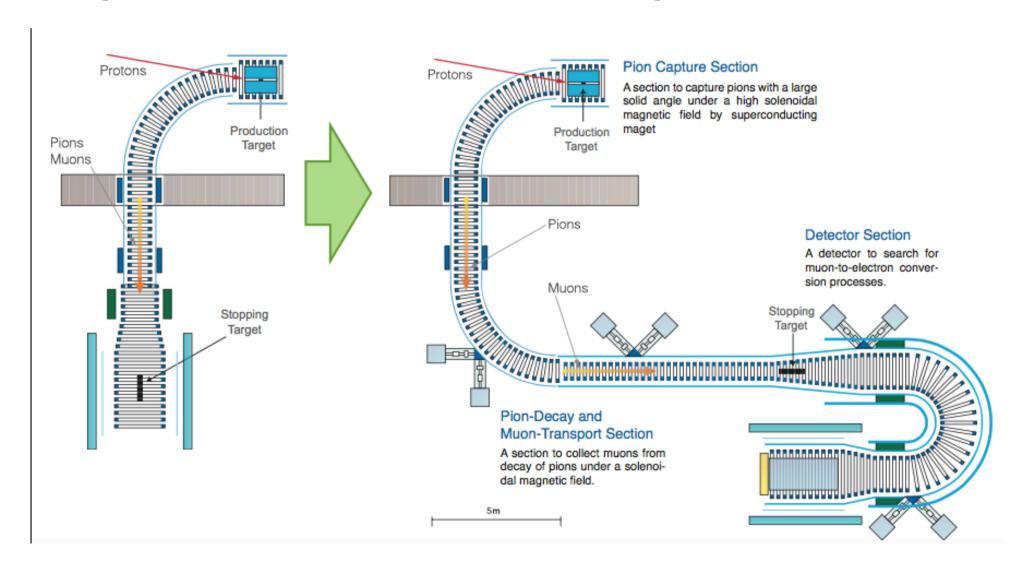


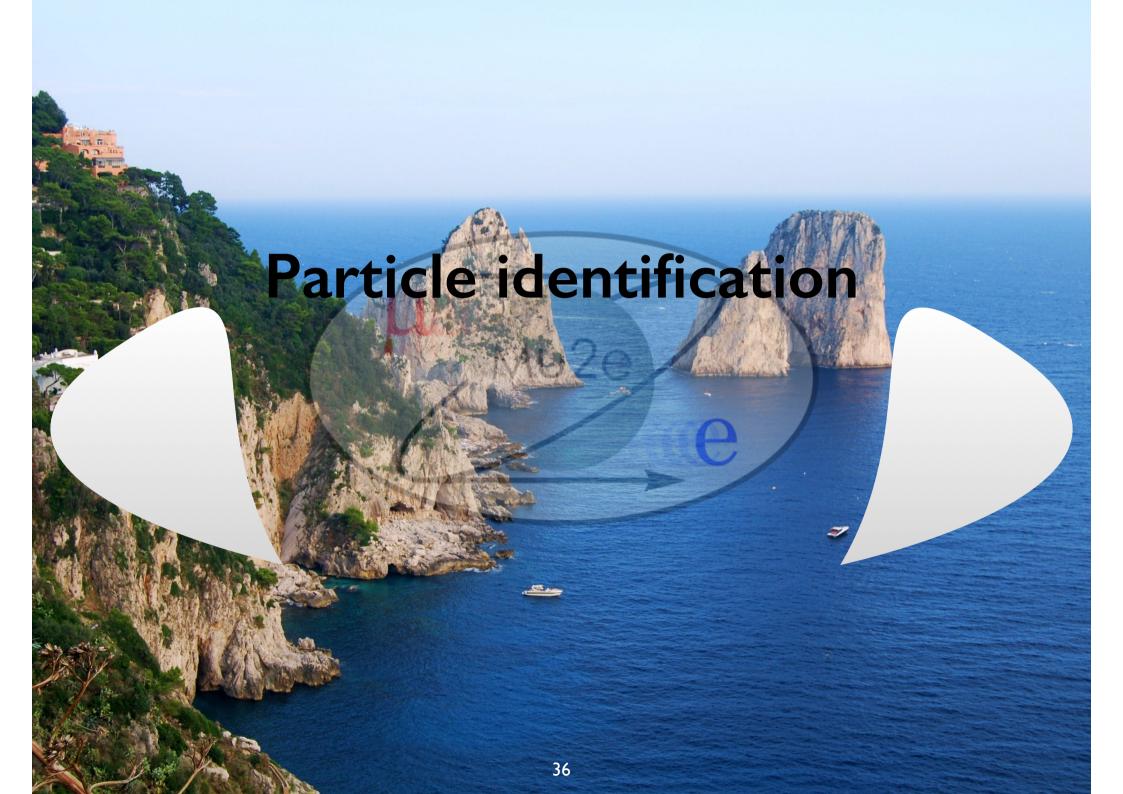
COMET experiment



phase I

phase II





Why Particle Identification is needed



- Cosmic ray and antiproton induced background can be divided into 2 main categories:
 - I. e generated via interactions producing a track mimicking the CE
 - 2. non-electron particles (μ and π) that are reconstructed as an "electron-like" track mimicking the CE
- (I) represents the irreducible background, while (2) can be suppressed using a PID method

Mu2e PID method:

- Information from reconstructed tracks and calorimeter clusters are combined for identifying group (2)
- Stringent requirement from Cosmic: μ-rejection factor ≥ 200

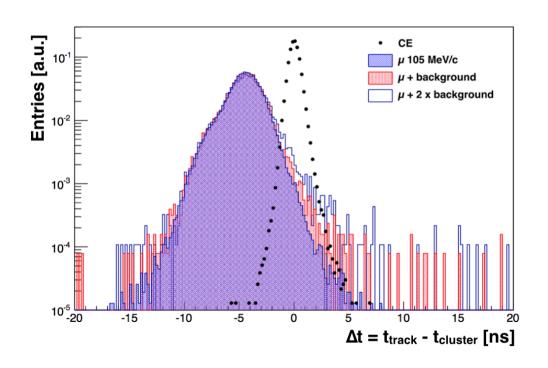


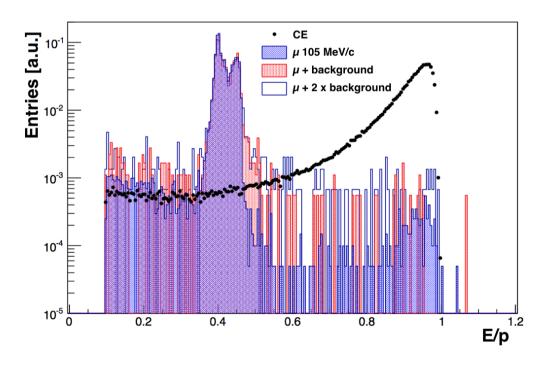
Cosmic µ rejection



- 105 MeV/c e⁻ are ultra-relativistic, while 105 MeV/c μ have β ~ 0.7 and a kinetic energy of ~ 40 MeV
- Likelihood rejection combines $\Delta t = t_{track} t_{cluster}$ and E/p:

$$\ln L_{e,\mu} = \ln P_{e,\mu}(\Delta t) + \ln P_{e,\mu}(E/p)$$

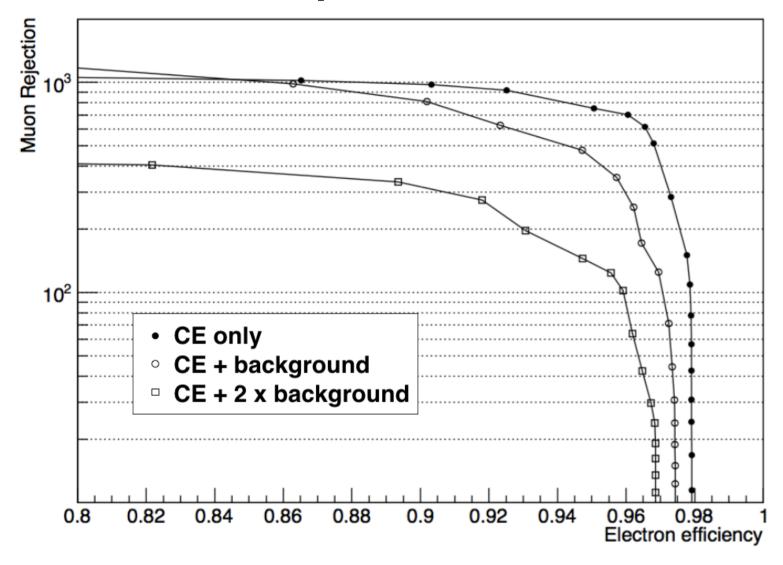






PID performance



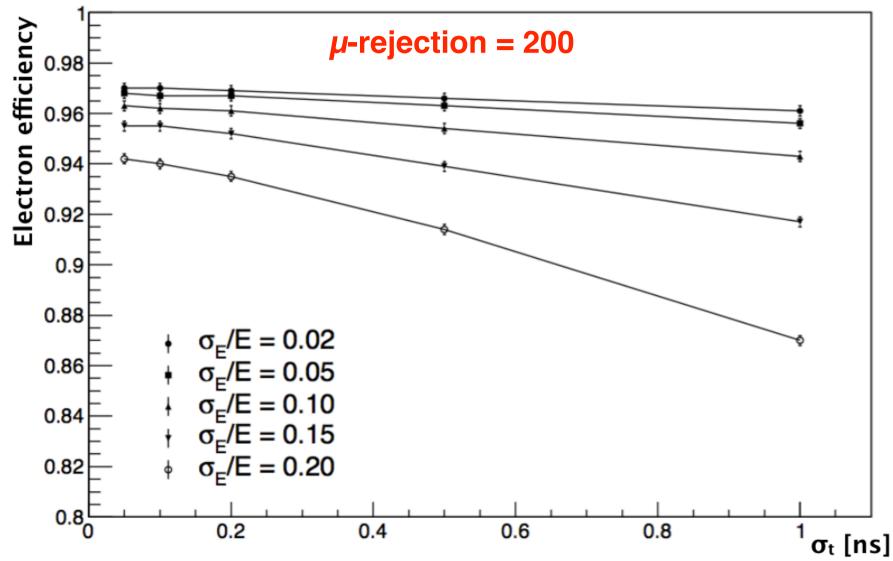


• A muon-rejection of 200 corresponds to a cut at ln $L_{e/\mu}$ > 1.5 and an e^- efficiency of ~ 96%



PID yield vs performance





• In the range $\sigma_E/E<0.1$ and $\sigma_t<0.5$ ns the e^- efficiency is within 2%





Crystals properties

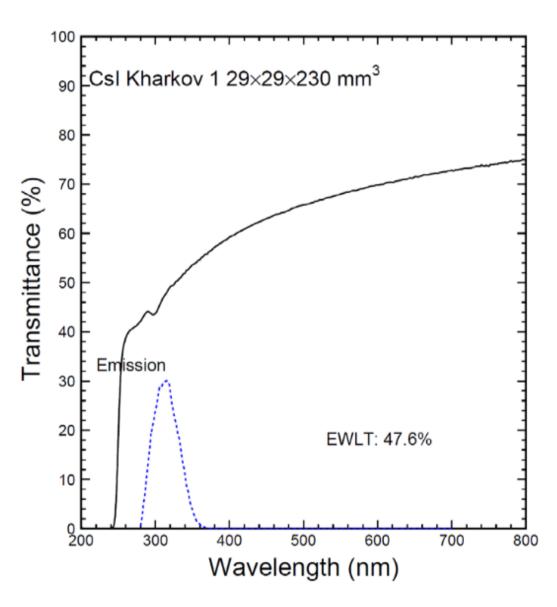


Crystal	$oxed{BaF_2}$	LYSO	CsI
Density [g/cm ³]	4.89	7.28	4.51
Radiation length [cm] X ₀	2.03	1.14	1.86
Molière radius [cm] R _m	3.10	2.07	3.57
Interaction length [cm]	30.7	20.9	39.3
dE/dx [MeV/cm]	6.5	10.0	5.56
Refractive Index at λ_{max}	1.50	1.82	1.95
Peak luminescence [nm]	220, 300	402	310
Decay time τ [ns]	0.9, 650	40	16
Light yield (compared to NaI(TI)) [%]	4.1, 3.6	85	3.6
Light yield variation with	0.1, -1.9	-0.2	-1.4
temperature $[\%/^{\circ}C]$			
Hygroscopicity	Slight	None	Slight



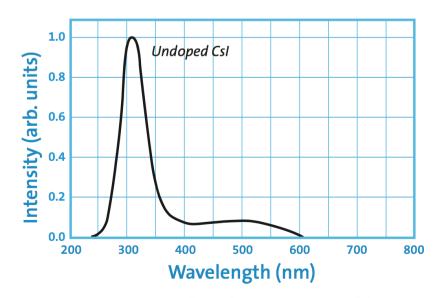
Csl properties





$$EWLT = \frac{\int LT(\lambda)Em(\lambda)d\lambda}{\int Em(\lambda)d\lambda}$$

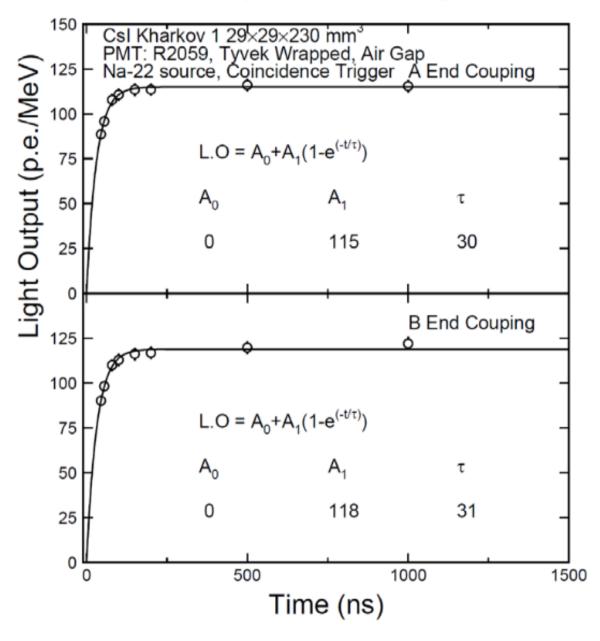
• where $LT(\lambda)$ is the light transmittance and $Em(\lambda)$ is the emission spectrum





Csl Light Output

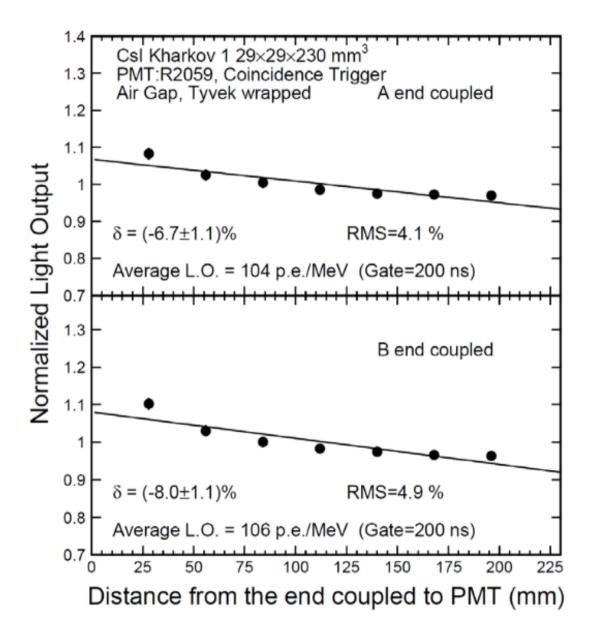






Csl LRU

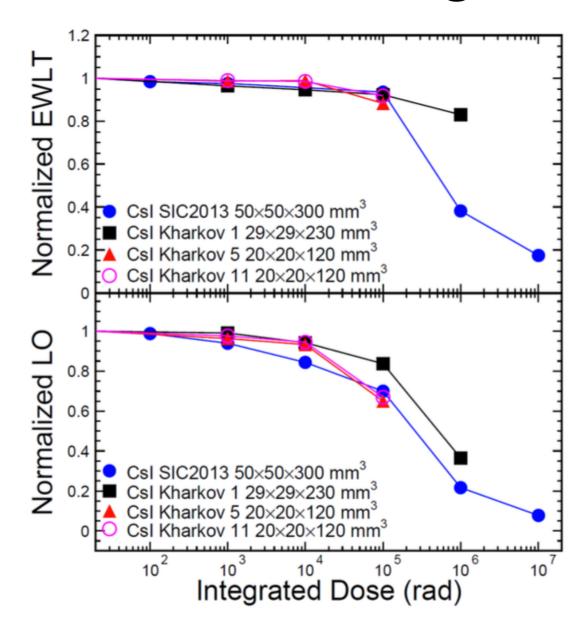






Csl rad damage

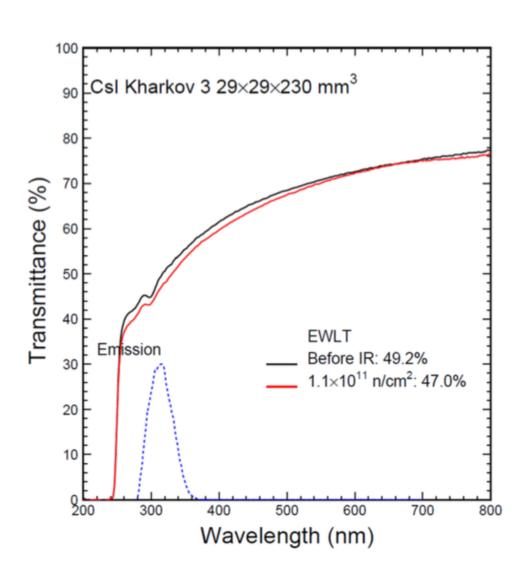


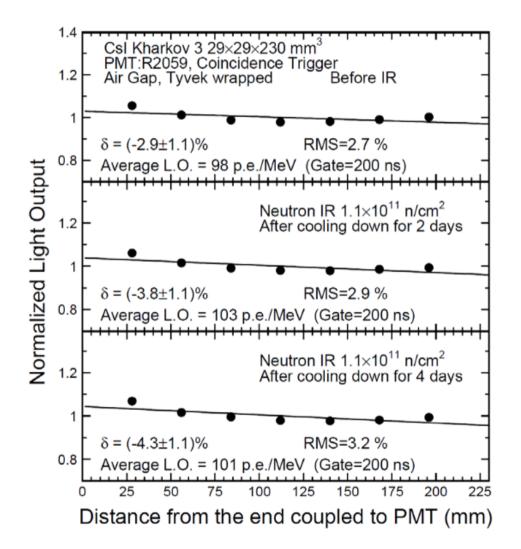




Csl neutron damage





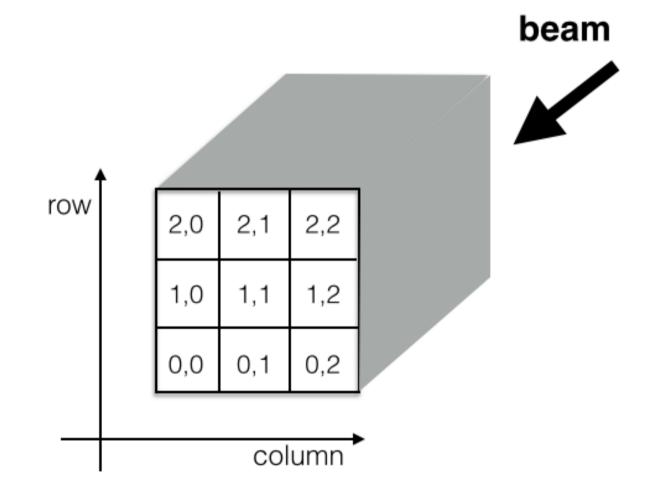






MUSE Numbering convention



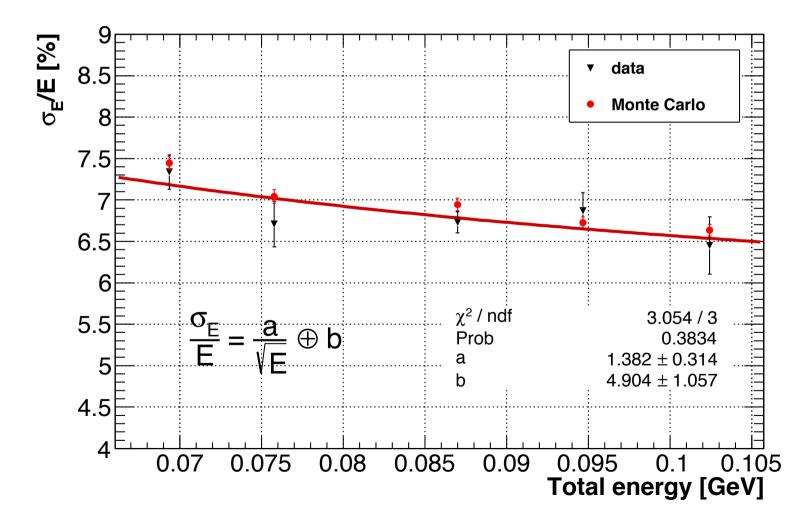




Energy resolution



- Prototype dimensions: I.3 R_{Moliere}² x I0 X₀
- Still comparison between data and Monte Carlo useful

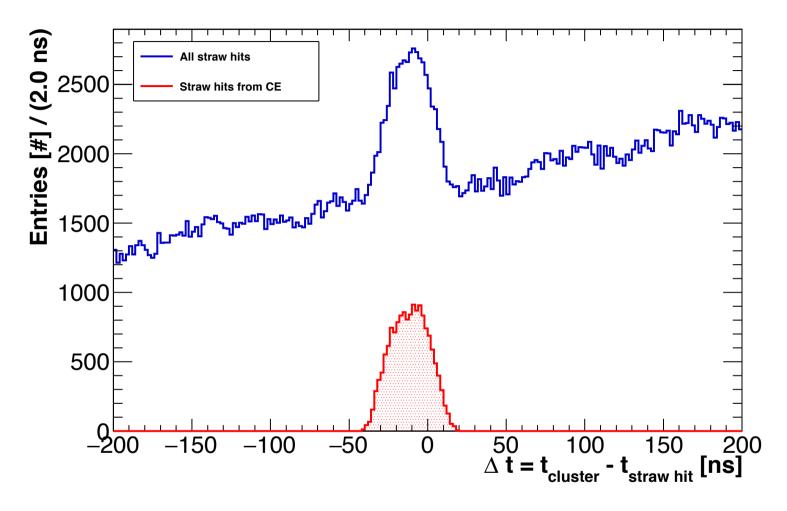






Cluster time window

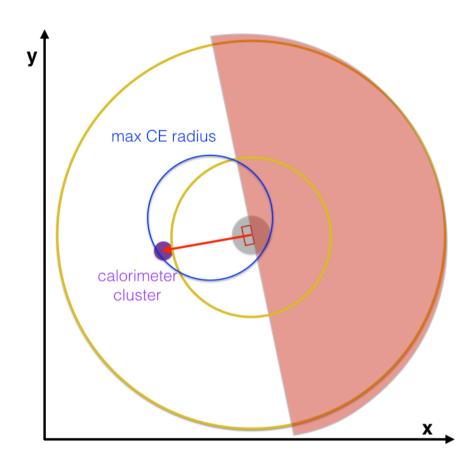




- Average tof from middle of the tracker to the calorimeter ~ 8 ns
- Mean drift time ~ 20 ns
- Difference of these two numbers is consistent with the peak position

MUSE Cluster positon selection





- Graded magnetic field between the stopping target and the tracker limits the CE pT
- Cluster position identifies the semi-plane where the CE track relies

MUSEMu2e track reconstruction



The Mu2e track reconstruction has several specific features:

- a CE makes 2-3 full turns in the tracker
- time dependence of the track-hit position:

$$r_{\text{drift}} = v_{\text{drift}} \cdot (t_{\text{measured}} - T_0 - t_{\text{flight}})$$

The track reconstruction is factorized into 2 main steps:

- 1. Track finding: provides a set of straw hits consistent with a track candidate
- 2. Kalman based fitter: performs the final reconstruction

The track finding uses two algorithms:

- A. Standalone: relies only on the tracker information
- B. Calorimeter-seeded: seeds the track search using the reco cluster